



## THEORY GUIDE

# Equations of Fluid Flow

## Kinetic Energy Equation in Cartesian and Cylindrical Coordinates

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## 1 Kinetic energy equation in Cartesian coordinates

In Ref. [1] we derived the continuity equation for a three-dimensional unsteady compressible fluid flow in Cartesian coordinates:

$$\frac{\partial \rho}{\partial t} + \frac{\partial(\rho u)}{\partial x} + \frac{\partial(\rho v)}{\partial y} + \frac{\partial(\rho w)}{\partial z} = 0 \quad [\text{kg m}^{-3} \text{ s}^{-1}] \quad (1.1)$$

In Ref. [2] we derived the equations for the three components of momentum in a three-dimensional unsteady compressible fluid flow in Cartesian coordinates. When written in terms of stress, the equations are

$$\frac{\partial(\rho u)}{\partial t} + \frac{\partial(\rho u^2)}{\partial x} + \frac{\partial(\rho uv)}{\partial y} + \frac{\partial(\rho uw)}{\partial z} = -\frac{\partial p}{\partial x} + \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} + \rho f_x \quad [\text{kg m}^{-2} \text{ s}^{-2}] \quad (1.2)$$

$$\frac{\partial(\rho v)}{\partial t} + \frac{\partial(\rho vu)}{\partial x} + \frac{\partial(\rho v^2)}{\partial y} + \frac{\partial(\rho vw)}{\partial z} = -\frac{\partial p}{\partial y} + \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} + \rho f_y \quad [\text{kg m}^{-2} \text{ s}^{-2}] \quad (1.3)$$

$$\frac{\partial(\rho w)}{\partial t} + \frac{\partial(\rho wu)}{\partial x} + \frac{\partial(\rho wv)}{\partial y} + \frac{\partial(\rho w^2)}{\partial z} = -\frac{\partial p}{\partial z} + \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} + \rho f_z \quad [\text{kg m}^{-2} \text{ s}^{-2}] \quad (1.4)$$

The kinetic energy per unit mass of the fluid flow  $K$  is  $(u^2 + v^2 + w^2)/2$ . We can derive an equation for  $K$  from the continuity equation and the three momentum equations. The left-hand side of the  $x$ -component momentum equation (1.2) can be expanded to give

$$\begin{aligned} \rho \frac{\partial u}{\partial t} + \rho u \frac{\partial u}{\partial x} + \rho v \frac{\partial u}{\partial y} + \rho w \frac{\partial u}{\partial z} + u \frac{\partial \rho}{\partial t} + u \frac{\partial(\rho u)}{\partial x} + u \frac{\partial(\rho v)}{\partial y} + u \frac{\partial(\rho w)}{\partial z} \\ = -\frac{\partial p}{\partial x} + \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} + \rho f_x \end{aligned}$$

From the continuity equation (1.1) the last four terms on the left-hand side of this equation are zero, so

$$\rho \frac{\partial u}{\partial t} + \rho u \frac{\partial u}{\partial x} + \rho v \frac{\partial u}{\partial y} + \rho w \frac{\partial u}{\partial z} = -\frac{\partial p}{\partial x} + \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} + \rho f_x$$

By multiplying through by  $u$ , we have

$$\rho \frac{\partial \left(\frac{u^2}{2}\right)}{\partial t} + \rho u \frac{\partial \left(\frac{u^2}{2}\right)}{\partial x} + \rho v \frac{\partial \left(\frac{u^2}{2}\right)}{\partial y} + \rho w \frac{\partial \left(\frac{u^2}{2}\right)}{\partial z} = -u \frac{\partial p}{\partial x} + u \left[ \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \right] + \rho u f_x \quad (1.5)$$

Similarly, for the  $y$ - and  $z$ -component momentum equations:

$$\rho \frac{\partial \left(\frac{v^2}{2}\right)}{\partial t} + \rho u \frac{\partial \left(\frac{v^2}{2}\right)}{\partial x} + \rho v \frac{\partial \left(\frac{v^2}{2}\right)}{\partial y} + \rho w \frac{\partial \left(\frac{v^2}{2}\right)}{\partial z} = -v \frac{\partial p}{\partial y} + v \left[ \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} \right] + \rho v f_y \quad (1.6)$$

$$\rho \frac{\partial \left(\frac{w^2}{2}\right)}{\partial t} + \rho u \frac{\partial \left(\frac{w^2}{2}\right)}{\partial x} + \rho v \frac{\partial \left(\frac{w^2}{2}\right)}{\partial y} + \rho w \frac{\partial \left(\frac{w^2}{2}\right)}{\partial z} = -w \frac{\partial p}{\partial z} + w \left[ \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right] + \rho w f_z \quad (1.7)$$

By summing (1.5), (1.6) and (1.7), we obtain

$$\begin{aligned} & \rho \frac{\partial K}{\partial t} + \rho u \frac{\partial K}{\partial x} + \rho v \frac{\partial K}{\partial y} + \rho w \frac{\partial K}{\partial z} \\ &= - \left[ u \frac{\partial p}{\partial x} + v \frac{\partial p}{\partial y} + w \frac{\partial p}{\partial z} \right] \\ &+ u \left[ \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \right] + v \left[ \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} \right] + w \left[ \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right] \\ & \rho u f_x + \rho v f_y + \rho w f_z \quad [\text{J m}^{-3} \text{ s}^{-1}] \quad (1.8) \end{aligned}$$

Multiplying the continuity equation (1.1) by  $K$  gives

$$K \frac{\partial \rho}{\partial t} + K \frac{\partial(\rho u)}{\partial x} + K \frac{\partial(\rho v)}{\partial y} + K \frac{\partial(\rho w)}{\partial z} = 0$$

and adding this equation to the left-hand side of (1.8) gives

$$\begin{aligned} & \frac{\partial(\rho K)}{\partial t} + \frac{\partial(\rho u K)}{\partial x} + \frac{\partial(\rho v K)}{\partial y} + \frac{\partial(\rho w K)}{\partial z} \\ &= - \left[ u \frac{\partial p}{\partial x} + v \frac{\partial p}{\partial y} + w \frac{\partial p}{\partial z} \right] \\ &+ u \left[ \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{xy}}{\partial y} + \frac{\partial \sigma_{xz}}{\partial z} \right] + v \left[ \frac{\partial \sigma_{yx}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{yz}}{\partial z} \right] + w \left[ \frac{\partial \sigma_{zx}}{\partial x} + \frac{\partial \sigma_{zy}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right] \\ & \rho u f_x + \rho v f_y + \rho w f_z \quad [\text{J m}^{-3} \text{ s}^{-1}] \quad (1.9) \end{aligned}$$

which is the equation for the conservation of kinetic energy  $K$  in Cartesian coordinates.

## 2 Kinetic energy equation in cylindrical coordinates

In Ref. [1] we derived the continuity equation for a three-dimensional unsteady compressible fluid flow in cylindrical coordinates:

$$\frac{\partial \rho}{\partial t} + \frac{1}{r} \frac{\partial(r\rho v_r)}{\partial r} + \frac{1}{r} \frac{\partial(\rho v_\theta)}{\partial \theta} + \frac{\partial(\rho v_z)}{\partial z} = 0 \quad [\text{kg m}^{-3} \text{ s}^{-1}] \quad (2.1)$$

In Ref. [3] we derived the equations for the three components of momentum in a three-dimensional unsteady compressible fluid flow in cylindrical coordinates. When written in terms of stress, the equations are

$$\begin{aligned} & \frac{\partial(\rho v_r)}{\partial t} + \frac{1}{r} \frac{\partial(r\rho v_r^2)}{\partial r} + \frac{1}{r} \frac{\partial(\rho v_r v_\theta)}{\partial \theta} + \frac{\partial(\rho v_r v_z)}{\partial z} \\ &= -\frac{\partial p}{\partial r} + \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} + \rho f_r \end{aligned} \quad (2.2)$$

$$\begin{aligned} & \frac{\partial(\rho v_\theta)}{\partial t} + \frac{1}{r} \frac{\partial(r\rho v_r v_\theta)}{\partial r} + \frac{1}{r} \frac{\partial(\rho v_\theta^2)}{\partial \theta} + \frac{\partial(\rho v_\theta v_z)}{\partial z} \\ &= -\frac{1}{r} \frac{\partial p}{\partial \theta} + \frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\sigma_{\theta r}}{r} + \frac{\partial \sigma_{\theta r}}{\partial r} + \frac{\partial \sigma_{\theta z}}{\partial z} + \frac{\sigma_{r\theta}}{r} + \rho f_\theta \end{aligned} \quad (2.3)$$

$$\begin{aligned} & \frac{\partial(\rho v_z)}{\partial t} + \frac{1}{r} \frac{\partial(r\rho v_r v_z)}{\partial r} + \frac{1}{r} \frac{\partial(\rho v_\theta v_z)}{\partial \theta} + \frac{\partial(\rho v_z^2)}{\partial z} \\ &= -\frac{\partial p}{\partial z} + \frac{1}{r} \frac{\partial \sigma_{zr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{z\theta}}{\partial \theta} + \frac{\partial \sigma_{zz}}{\partial z} + \rho f_z \end{aligned} \quad (2.4)$$

The kinetic energy per unit mass of the fluid flow  $K$  is  $(v_r^2 + v_\theta^2 + v_z^2)/2$ . We can derive an equation for  $K$  from the continuity equation and the three momentum equations. The left-hand side of the  $r$ -component momentum equation (2.2) can be expanded to give

$$\begin{aligned} & \rho \frac{\partial v_r}{\partial t} + \rho v_r \frac{\partial v_r}{\partial r} + \frac{\rho v_\theta}{r} \frac{\partial v_r}{\partial \theta} + \rho v_z \frac{\partial v_r}{\partial z} + v_r \frac{\partial \rho}{\partial t} + \frac{v_r}{r} \frac{\partial(r\rho v_r)}{\partial r} + \frac{v_r}{r} \frac{\partial(\rho v_\theta)}{\partial \theta} + v_r \frac{\partial(\rho v_z)}{\partial z} \\ &= -\frac{\partial p}{\partial r} + \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} + \rho f_r \end{aligned}$$

From the continuity equation (2.1) the last four terms on the left-hand side of this equation are zero, so

$$\begin{aligned} & \rho \frac{\partial v_r}{\partial t} + \rho v_r \frac{\partial v_r}{\partial r} + \frac{\rho v_\theta}{r} \frac{\partial v_r}{\partial \theta} + \rho v_z \frac{\partial v_r}{\partial z} \\ &= -\frac{\partial p}{\partial r} + \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} + \rho f_r \end{aligned}$$

By multiplying through by  $v_r$ , we have

$$\begin{aligned} & \rho \frac{\partial \left( \frac{v_r^2}{2} \right)}{\partial t} + \rho v_r \frac{\partial \left( \frac{v_r^2}{2} \right)}{\partial r} + \frac{\rho v_\theta}{r} \frac{\partial \left( \frac{v_r^2}{2} \right)}{\partial \theta} + \rho v_z \frac{\partial \left( \frac{v_r^2}{2} \right)}{\partial z} \\ &= -v_r \frac{\partial p}{\partial r} + v_r \left[ \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} \right] + \rho v_r f_r \end{aligned} \quad (2.5)$$

Similarly, for the  $\theta$ - and  $z$ -component momentum equations:

$$\begin{aligned} & \rho \frac{\partial \left( \frac{v_\theta^2}{2} \right)}{\partial t} + \rho v_r \frac{\partial \left( \frac{v_\theta^2}{2} \right)}{\partial x} + \frac{\rho v_\theta}{r} \frac{\partial \left( \frac{v_\theta^2}{2} \right)}{\partial y} + \rho v_z \frac{\partial \left( \frac{v_\theta^2}{2} \right)}{\partial z} \\ &= -\frac{v_\theta}{r} \frac{\partial p}{\partial \theta} + v_\theta \left[ \frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\sigma_{\theta r}}{r} + \frac{\partial \sigma_{\theta r}}{\partial r} + \frac{\partial \sigma_{\theta z}}{\partial z} + \frac{\sigma_{r\theta}}{r} \right] + \rho v_\theta f_\theta \end{aligned} \quad (2.6)$$

$$\begin{aligned} & \rho \frac{\partial \left( \frac{v_z^2}{2} \right)}{\partial t} + \rho v_r \frac{\partial \left( \frac{v_z^2}{2} \right)}{\partial x} + \frac{\rho v_\theta}{r} \frac{\partial \left( \frac{v_z^2}{2} \right)}{\partial y} + \rho v_z \frac{\partial \left( \frac{v_z^2}{2} \right)}{\partial z} \\ &= -v_z \frac{\partial p}{\partial z} + v_z \left[ \frac{1}{r} \frac{\partial r \sigma_{zr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{z\theta}}{\partial \theta} + \frac{\partial \sigma_{zz}}{\partial z} \right] + \rho v_z f_z \end{aligned} \quad (2.7)$$

By summing (2.5), (2.6) and (2.7), we obtain

$$\begin{aligned} & \rho \frac{\partial K}{\partial t} + \rho v_r \frac{\partial K}{\partial r} + \frac{\rho v_\theta}{r} \frac{\partial K}{\partial \theta} + \rho v_z \frac{\partial K}{\partial z} \\ &= - \left[ v_r \frac{\partial p}{\partial r} + \frac{v_\theta}{r} \frac{\partial p}{\partial \theta} + v_z \frac{\partial p}{\partial z} \right] \\ &+ v_r \left[ \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} \right] + v_\theta \left[ \frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\sigma_{\theta r}}{r} + \frac{\partial \sigma_{\theta r}}{\partial r} + \frac{\partial \sigma_{\theta z}}{\partial z} + \frac{\sigma_{r\theta}}{r} \right] \\ &+ v_z \left[ \frac{1}{r} \frac{\partial r \sigma_{zr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{z\theta}}{\partial \theta} + \frac{\partial \sigma_{zz}}{\partial z} \right] \\ & \rho v_r f_r + \rho v_\theta f_\theta + \rho v_z f_z \quad [\text{J kg}^{-1} \text{ s}^{-1}] \end{aligned} \quad (2.8)$$



Multiplying the continuity equation (2.1) by  $K$  gives

$$K \frac{\partial \rho}{\partial t} + \frac{K}{r} \frac{\partial (r \rho v_r)}{\partial r} + \frac{K}{r} \frac{\partial (\rho v_\theta)}{\partial \theta} + K \frac{\partial (\rho v_z)}{\partial z} = 0$$

and adding this equation to the left-hand side of (2.8) gives

$$\begin{aligned} & \frac{\partial \rho K}{\partial t} + \frac{1}{r} \frac{\partial (r \rho v_r K)}{\partial r} + \frac{1}{r} \frac{\partial (\rho v_\theta K)}{\partial \theta} + \rho v_z \frac{\partial (\rho v_z K)}{\partial z} \\ &= - \left[ v_r \frac{\partial p}{\partial r} + \frac{v_\theta}{r} \frac{\partial p}{\partial \theta} + v_z \frac{\partial p}{\partial z} \right] \\ &+ v_r \left[ \frac{\partial \sigma_{rr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{r\theta}}{\partial \theta} + \frac{\partial \sigma_{rz}}{\partial z} + \frac{\sigma_{rr}}{r} - \frac{\sigma_{\theta\theta}}{r} \right] + v_\theta \left[ \frac{1}{r} \frac{\partial \sigma_{\theta\theta}}{\partial \theta} + \frac{\sigma_{\theta r}}{r} + \frac{\partial \sigma_{\theta r}}{\partial r} + \frac{\partial \sigma_{\theta z}}{\partial z} + \frac{\sigma_{r\theta}}{r} \right] \\ &+ v_z \left[ \frac{1}{r} \frac{\partial r \sigma_{zr}}{\partial r} + \frac{1}{r} \frac{\partial \sigma_{z\theta}}{\partial \theta} + \frac{\partial \sigma_{zz}}{\partial z} \right] \\ &\rho v_r f_r + \rho v_\theta f_\theta + \rho v_z f_z \quad [\text{J kg}^{-1} \text{s}^{-1}] \quad (2.9) \end{aligned}$$

which is the equation for the conservation of kinetic energy  $K$  in cylindrical coordinates.



### 3 References

The following reports may be downloaded from

<https://atkinsonscience.co.uk/Downloads/FluidDynamics.aspx>

1. K. N. Atkinson, *Equations of Fluid Flow, Continuity Equation in Cartesian and Cylindrical Coordinates, Theory Guide*, Atkinson Science Limited, 2020.
2. K. N. Atkinson, *Equations of Fluid Flow, Momentum Equations in Cartesian Coordinates, Theory Guide*, Atkinson Science Limited, 2020.
3. K. N. Atkinson, *Equations of Fluid Flow, Momentum Equations in Cylindrical Coordinates, Theory Guide*, Atkinson Science Limited, 2020.